

Quantum Homotopy

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Stephen Peña - Mechanics on Manifolds
and Classical Field Theory

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- Outline** -
- 1) Classical Mechanics
 - 2) Quantum Mechanics
 - 3) Gauge Theory
 - 4) Functorial Quantum Field Theory (FQFT), Factorization Algebras

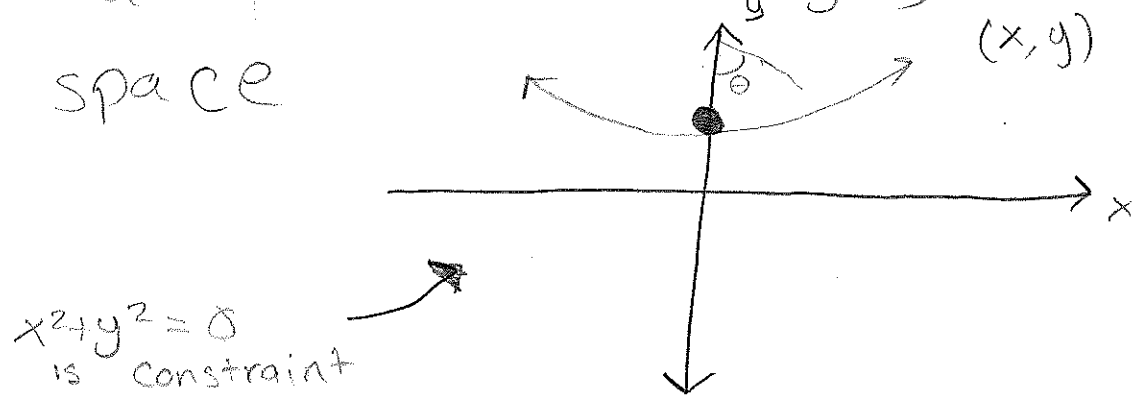
Mechanics is dominated by Hamilton's

Principle of Least Action: We

discuss generalized coordinates. Let's

look at a pendulum swinging

in x - y space



Def: Constraints are equations your system must obey.

Def: A motion in \mathbb{R}^n is a differentiable mapping $\vec{x}: I \rightarrow \mathbb{R}^n$

$\dot{\vec{x}}$ is short hand for $\frac{d\vec{x}}{dt}$. Also

$\dot{\vec{x}}$ = Velocity

$\ddot{\vec{x}}$ = acceleration

Def: A motion of a systems of n-points is given by n mappings $x_i: \mathbb{R} \rightarrow \mathbb{R}^N$.

Def: The direct Product of n copies of \mathbb{R}^N is called configuration space.

We have

$$F: \mathbb{R}^N \times \mathbb{R}^N \times \mathbb{R} \rightarrow \mathbb{R}^N$$

Such that $\ddot{\vec{x}} = F(x, \dot{x}, t)$ \leftarrow (F=ma)

Def: A system with one degree of freedom is a system defined by one differential Equation: (2)

$$\ddot{x} = f(x), \quad x \in \mathbb{R}$$

Def: The Kinetic energy is the quadratic form

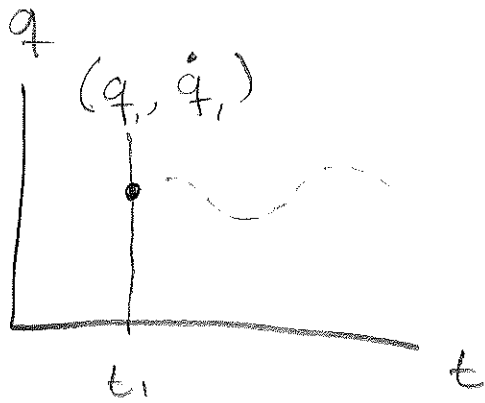
$$T = \frac{1}{2} \dot{x}^2$$

Def: The Potential energy is the function

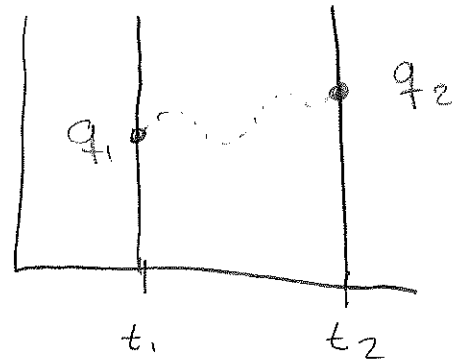
$$U(x) = - \int_{x_0}^x f(\xi) d\xi$$

The total energy is $E = T + U$.

Two ways to view Motion



IVP



"boundary" value

Hamilton's Principle: Motion of a system in a time interval $t \in (t_1, t_2)$ coincides with the extremum of the action functional.

$$S = S[q] = \int_{t_1}^{t_2} L(q, \dot{q}, t) dt$$

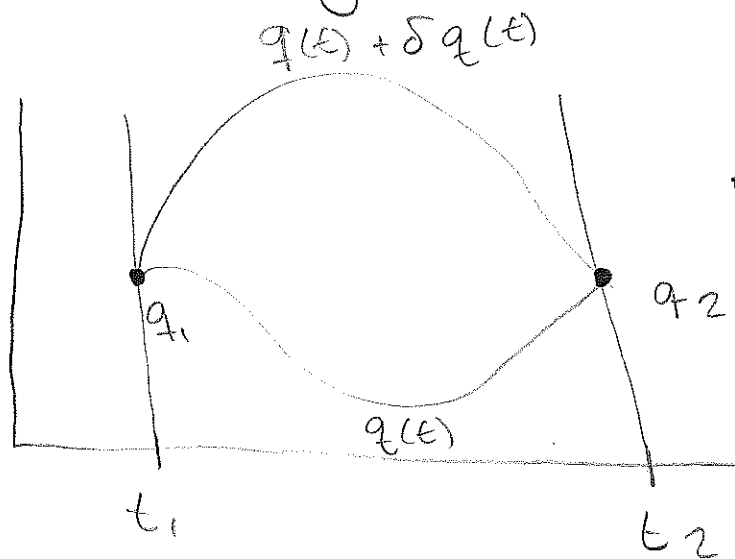
↑
Mind e: function

"Lagrangien"

How do we get it? Calculus of Variation.

Let $X = (\mathbb{R}^N)^I (= \{\text{Curves in Space}\})$ ③

Consider again



we will vary curves

$\delta q(t)$ vanishes at t_1, t_2 .

$$* \delta \frac{d}{dt} = \frac{d}{dt} \delta$$

So

$$\begin{aligned} \delta S &= \delta \int_{t_1}^{t_2} L dt \\ &= \int_{t_1}^{t_2} \delta L dt \end{aligned}$$

$$= \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q} \delta q + \frac{\partial L}{\partial \dot{q}} \delta \dot{q} \right) dt$$

$$= \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q} \delta q + \frac{\partial L}{\partial \dot{q}} \frac{d}{dt} \delta q \right) dt$$

$$= \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q} \delta q + \frac{d}{dt} \frac{\partial L}{\partial \dot{q}} \delta q \right) dt$$

$$\stackrel{\text{by parts}}{=} \int_{t_1}^{t_2} \left(\frac{\partial L}{\partial q} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}} \right) \right) \delta q dt + \left[\frac{\partial L}{\partial \dot{q}} \delta q \right]_{t_1}^{t_2}$$

= 0

Since $\delta q \neq 0$ we have

$$\frac{\partial L}{\partial q} - \frac{d}{dt} \left(\frac{\partial L}{\partial \dot{q}} \right) = 0 \quad (*)$$

↗
Euler-Lagrange Equations

Remarks: 1) 2nd order ODE

(4)

$$2) L = T - V$$

$$3) L'(q, \dot{q}, t) = L + \frac{d\lambda(q, t)}{dt}$$

To solve (*) we have the following ways

i) Analytically (Integrable Systems)

ii) Perturbations

iii) Numerical

Hamiltonian Mechanics:

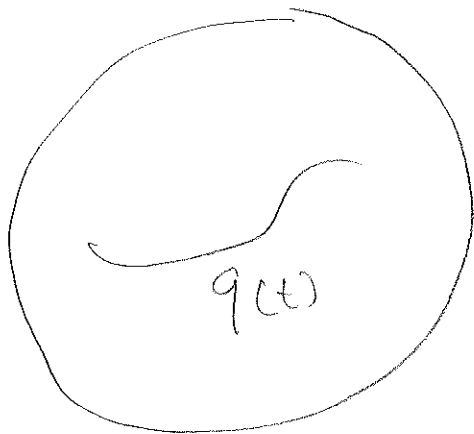
Why another formalism?

Trade E-L \longrightarrow Hamilton's Equations

$n - 2^{\text{nd}}$ order

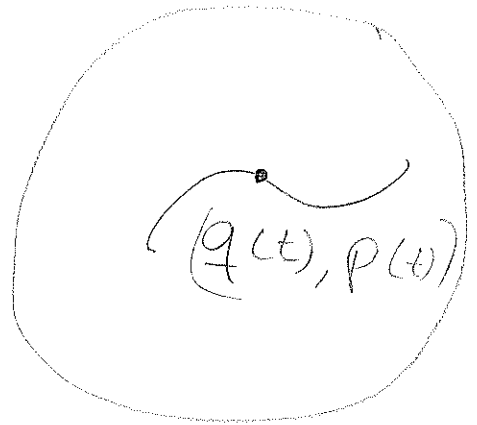
$2n - 1^{\text{st}}$ order

LM



$(t, q(t))$

HM Phase space



$\{q^i, p_i\}$

where
$$p_i := \frac{\partial L}{\partial \dot{q}^i}$$

Hamiltonian = Legendre transform
of Lagrangian

$$H(q, p, t) = p_i \dot{q}^i - L$$

Hamilton's Theorem: Systems of
n E-L is equivalent to 2n 1st
order hamiltas equations

$$p_j = - \frac{\partial H}{\partial q^j}$$

$$\dot{q}^j = \frac{\partial H}{\partial p_j}$$

Similar process to E-2 we (5)
have

$$0 = \delta S = \int_{t_1}^{t_2} \left[\delta p \left(\dot{q} - \frac{\partial H}{\partial p} \right) + \delta q \left(-\frac{\partial H}{\partial q} - \dot{p} \right) \right] dt$$

Def: Let f, g be functions on Phase space, then their poisson bracket is a new phase space function

$$\{f, g\} = \sum_{p_i, q_i} \left(\frac{\partial f}{\partial q_i} \frac{\partial g}{\partial p_i} - \frac{\partial f}{\partial p_i} \frac{\partial g}{\partial q_i} \right)$$

$$\{q^i, q^j\} = 0 = \{p_i, p_j\}$$

$$\{q^i, p_j\} = \delta_{ij}$$

